



## Can The Spectra Of Hermitian Operator Be Invariant

### Under $x \leftrightarrow p$ ? :Case Study :1D General Oscillator

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#### Abstract

We address an intriguing question on spectral invariance in quantum mechanics on exchange of co-ordinate and momentum  $x \leftrightarrow p$  considering general oscillator as an example.

#### Keywords

exchange of co-ordinate and momentum , spectral invariance , general oscillator .

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#### I.Introduction

In a recent paper Rath and Mallick [1] proposed a generalised model on co-ordinate and momentum transformation in the case of Harmonic Oscillator to reflect the spectral invariance. Further it is well known that commutation relation

$$[x, p] = i \quad (1)$$

between co-ordinate ( $x$ ) and momentum ( $p$ ) on exchange( $x \leftrightarrow p$ ) becomes

$$[p, x] = -i \quad (2)$$

.Now question arises whether spectra of Hermitian operator (more precisely self-adjoint operator) be invariant under exchange of co-ordinate and momentum ?. If the answer to this case is yes , then why not address this to some model Hermitian operator. . In this context we would like to state that in the past there was a considerable interest among many others to study spectra of anharmonic oscillator [2-9] . In any way we consider a more general type of oscillator [2-9]

$$h = \mu p^2 + \lambda_1 x^2 + \lambda_2 x^4 + \lambda_3 x^6 \quad (3)$$

and study its spectra on exchange of co-ordinate and momentum.

#### II.New Operator and Spectra

Here we consider the operator

$$H = \mu x^2 + \lambda_1 p^2 + \lambda_2 p^4 + \lambda_3 p^6 \quad (4)$$

In order to solve it we use the eigenvalue relation[2,6,8,10,11]

$$H|\Psi\rangle = E|\Psi\rangle \quad (5)$$



where

$$|\psi\rangle = \sum A_m |m\rangle \quad (6)$$

In the above  $|m\rangle$  stands for standard harmonic oscillator wave function [10,11] satisfying the relation

$$[p^2 + x^2]|m\rangle = (2m + 1)|m\rangle \quad (7)$$

Now using the above relation one will notice that  $A_m$  satisfies the following recurrence relation

$$P_m A_{m-6} + Q_m A_{m-4} + R_m A_m + S_m A_m + T_m A_{m+2} + U_m A_{m+4} + V_m A_{m+6} = 0 \quad (8)$$

where

$$P_m = \langle m-6 | H | m \rangle \quad (9)$$

$$Q_m = \langle m-4 | H | m \rangle \quad (10)$$

$$R_m = \langle m-2 | H | m \rangle \quad (11)$$

$$S_m = \langle m | H | m \rangle \quad (12)$$

$$T_m = R_{m+2} \quad (13)$$

$$U_m = Q_{m+4} \quad (14)$$

$$V_m = P_{m+4} \quad (15)$$

The eigen values calculated using this relation using matrix diagonalisation method [8,10,11] are tabulated in table-1.



**Tale-1: Eigenvalues of New Operator and Comparision.**

n	$H = x^2 - 100p^2 + p^4$	$h = p^2 - 100x^2 + x^4$ [7]
0	- 2485.867 880 343	-2485.867 880 343
1	- 2485.867 880 343	-2485.867 880 343
2	-2457.643 822 699	-2457.643 822 699
3	-2457.643 822 699	-2457.643 822 699
n	$H = x^2 - 2p^2 - 2p^4 + p^6$	$h = p^2 - 2x^2 - 2x^4 + x^6$ [2,8]
0	-1.000 000	-0.999 987
1	-0.154 110	-0.154 093
2	3.629 625	3.629 880
3	8.007 560	8.007 742
n	$H = x^2 + p^4$	$h = p^2 + x^4$ [9,8]
0	1.060 362 090	1.060 362 090
1	3.799 073 029	3.799 073 029
2	7.455 697 937	7.455 697 973
3	11.644 745 511	11.644 745 511

### III.Conclusion

In one-dimensional general oscillator considered above we notice that hermitian operator has an equivalent operator whose eigenspectra remain invariant . Further we plot the  $|\Psi_{N=0-3}|^2$  corresponding to Hamiltonian

$$H = x^2 - 2p^2 - 2p^4 + p^6 \quad (16)$$

in fig-1 . Similarly we plot the  $|\Phi_{N=0-3}|^2$  corresponding to Hamiltonian

$$h = p^2 - 2x^2 - 2x^4 + x^6 \quad (17)$$

in fig-2. From the figs it is claar that eventhough two systems are iso-spectral in nature but different from each other.

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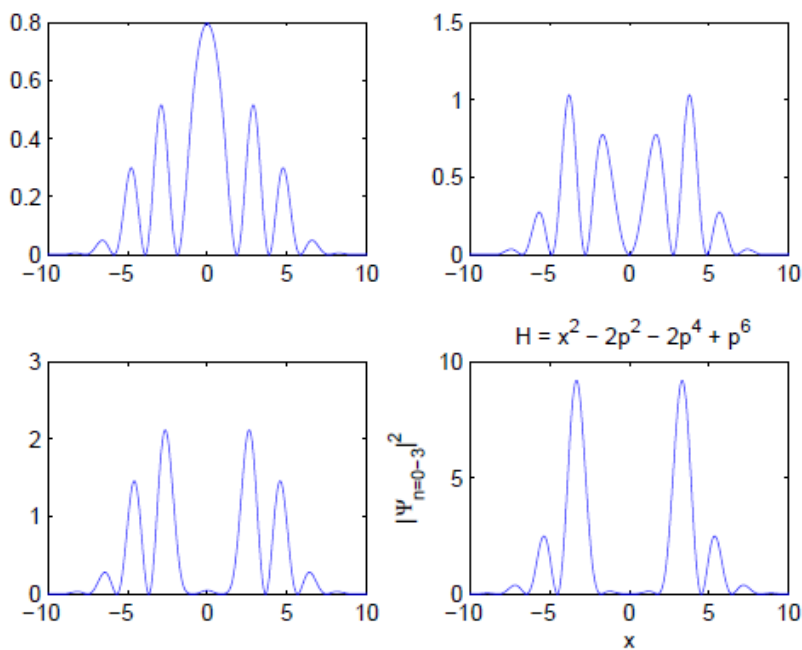


Figure 1:  $H = x^2 - 2p^2 - 2x^4 + x^6$   
 :  $|\Psi_{n=0-3}|^2$  of Equivalent Sextic Oscillator

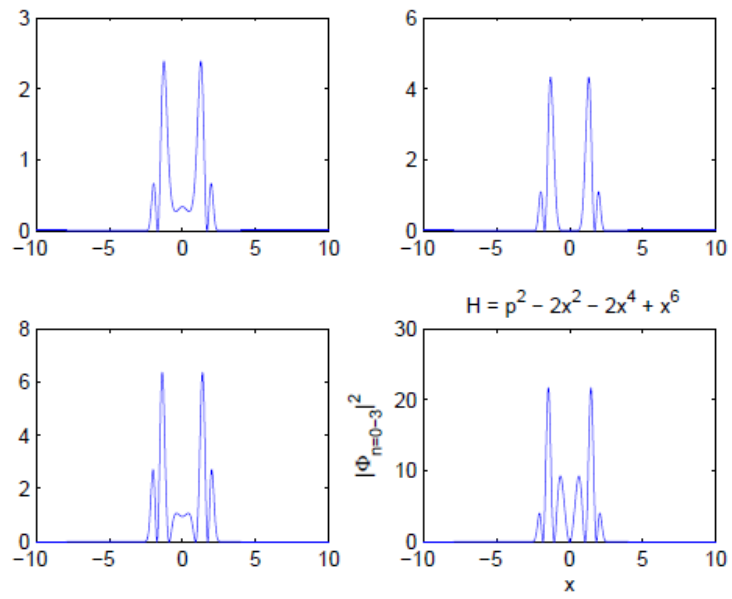


Figure 2:  $h = p^2 - 2x^2 - 2x^4 + x^6$   
:  $|\Phi_{n=0-3}|^2$  of Sextic Well Oscillator