



A Note on non-relativistic gravitation

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Abstract

Using the Jefimenko theory of gravitation [1] and the Le-Bellac, Levy-Leblond non relati-vistic approximations of Maxwell equations, we prove that in addition to newtonian gravi-tation, there exists a second non-relativistic gravitation.



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INTRODUCTION

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1. JEFIMENKO THEORY

The Jefimenko theory [1] of gravitation is obtained from electromagnetism by simply repla-cing the electromagnetic symbols and constants by the corresponding symbols and constants of gravitation in accordance with the following tables :

Table (1a) Electromagnetism

q (charge)
ρ (volume charge density)
J (convection current)
E (electric field)
B (magnetic field)
ϵ_0 (permittivity)
μ_0 (permeability)
$\epsilon_0\mu_0 = c^{-2}$

Table (1b) Gravitation

m (mass)
ρ (volume mass density)
J (mass current)
g (gravitational field)
K (cogravitational field)
$-1/4\pi G$
$-4\pi G/c^2$
G (gravitational constant)

Then, using these tables, the Maxwell and gravitational equations are in vacuum :

Maxwell (2a)

$$\begin{aligned} \nabla \cdot \mathbf{E} &= \rho / \epsilon_0 \\ \nabla \cdot \mathbf{B} &= 0 \\ \nabla \wedge \mathbf{E} &= -\partial_t \mathbf{B} \\ \nabla \wedge \mathbf{B} &= \mu_0 \mathbf{J} + 1/c^2 \partial_t \mathbf{E} \end{aligned}$$

Gravitation (2b)

$$\begin{aligned} \nabla \cdot \mathbf{g} &= -4\pi G \rho \\ \nabla \cdot \mathbf{K} &= 0 \\ \nabla \wedge \mathbf{g} &= -\partial_t \mathbf{K} \\ \nabla \wedge \mathbf{K} &= -4\pi G \mathbf{J} / c^2 + 1/c^2 \partial_t \mathbf{g} \end{aligned}$$

The cogravitational field **K** is a special feature of the Jefimenko theorY.

2 NON RELATIVISTIC GRAVITATION

Now, it is known [2,3,4] that the Maxwell equations have two nonrelativistic limits respecti-

vely called electric (3a) and magnetic limits(4a) supplying according to the tables (1a,b) the non relativistic limits of the gravitational equations that we call gravilimit (3b) and cogra-vilimit (4b) corresponding respectively to (3a) and (4a) :

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Electric limit (3a) [3]

$$\begin{aligned} \nabla \cdot \mathbf{E}_e &= \rho_e / \epsilon_0 \\ \nabla \cdot \mathbf{B}_e &= 0 \\ \nabla \wedge \mathbf{E}_e &= 0 \\ \nabla \wedge \mathbf{B}_e &= -1/c^2 \partial_t \mathbf{E}_e \end{aligned}$$

Magnetic limit (4a) [3]

$$\begin{aligned} \nabla \cdot \mathbf{E}_m &= \rho_m / \epsilon_0 \\ \nabla \cdot \mathbf{B}_m &= 0 \\ \nabla \wedge \mathbf{E}_m &= -\partial_t \mathbf{B}_m \\ \nabla \wedge \mathbf{B}_m &= \mu_0 \mathbf{J}_m \end{aligned}$$

Gravilimit (3b)

$$\begin{aligned} \nabla \cdot \mathbf{g}_g &= -4\pi G \rho_g \\ \nabla \cdot \mathbf{K}_g &= 0 \\ \nabla \wedge \mathbf{g}_g &= 0 \\ \nabla \wedge \mathbf{K}_g &= 1/c^2 \partial_t \mathbf{g}_g \end{aligned}$$

Cogravilimit (4b)

$$\begin{aligned} \nabla \cdot \mathbf{g}_{cg} &= -4\pi G \rho_{cg} \\ \nabla \cdot \mathbf{K}_{cg} &= 0 \\ \nabla \wedge \mathbf{g}_{cg} &= -\partial_t \mathbf{K}_{cg} \\ \nabla \wedge \mathbf{K}_{cg} &= -4\pi G \mathbf{J}_{cg} / c^2 \end{aligned}$$

3.GRAVITATION AND POTENTIALS

The electromagnetic potentials **V**, **A** relating to fields in vacuum have the magnetic limit [3]

satisfying the coulomb gauge with a scalar potential V_m independent on time $\partial_t V_m = 0$

$$\mathbf{E}_m = -\nabla V_m - \partial_t \mathbf{A}_m \quad , \quad \mathbf{B}_m = \nabla \wedge \mathbf{A}_m \quad , \quad \nabla \cdot \mathbf{A}_m = 0 \quad (5a)$$



Similarly in the electric limit, the potentials satisfy the Lorentz gauge and $\partial_t \mathbf{A}_e = 0$

$$\mathbf{E}_e = -\nabla V_e, \quad \mathbf{B}_e = \nabla \wedge \mathbf{A}_e, \quad \nabla \cdot \mathbf{A}_e + 1/c^2 \partial_t V_e = 0 \quad (6a)$$

Then, changing \mathbf{E}, \mathbf{B} , according to the tables (1,a,b) into the gravitational and cogravitational fields \mathbf{g}, \mathbf{K} and the subscripts e,m into g, cg gives similarly the gravilimit ($\partial_t \mathbf{A}_g = 0$) and the cogravitlimit ($\partial_t V_g = 0$)

$$\mathbf{g}_g = -\nabla V_g, \quad \mathbf{K}_g = \nabla \wedge \mathbf{A}_g, \quad \nabla \cdot \mathbf{A}_g + 1/c^2 \partial_t V_g = 0 \quad (6b)$$

$$\mathbf{g}_{cg} = -\nabla V_{cg} - \partial_t \mathbf{A}_{cg}, \quad \mathbf{K}_{cg} = \nabla \wedge \mathbf{A}_{cg}, \quad \nabla \cdot \mathbf{A}_{cg} = 0 \quad (5b)$$

The gravitational potentials have the Poisson approximation [1]

$$\nabla^2 V_g = 4\pi G\rho, \quad \nabla^2 \mathbf{A}_{cg} = 4\pi G\mathbf{J}/c^2 \quad (7)$$

So, the gravilimit supplies the newtonian gravitation generalized by the existence of the co-gravitational field \mathbf{K} in accordance with the Heaviside suggestion [5].

Remark : It has been proved [6] that E_g may be expressed in terms of a vector potential and K_{cg} in terms of a scalar potential.

4. CONCLUSION

The generalization of the basic gravitational equation of General Relativity theory is based at a large extent on the existence of the cogravitational field. Many applications of this gene-ralized gravity may be found in the Jefimenko books [1,7] (in a curious work [8] the advance of Mercury perihelion is explained by the cogravity).

We have proved here that the non relativistic approximation of the Jefimenko theory sup-plies in addition to a generalization of newtonian gravity (reducing in fact to newtonian gravity since the term due to the cogravitational field is very small and negligible), a second

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non relativistic gravity. The interpretation os these approximations fall out of the scope of this pre esnt Note.

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