# Comparison of Variational and Fractional Variational Iteration Methods in Solving Time-Fractional Differential Equations <br> ${ }^{1}$ Eman Ali Hussain, ${ }^{2}$ Zainab Mohammed Alwan <br> ${ }^{1}$ Asst. Prof, Dr., Department of Mathematics, College of Science, University of AL- Mustansiriyah, Iraq <br> E-mail: dr_emansultan@yahoo.com <br> ${ }^{2}$ Asst. Lecturer .,Department of Mathematics, College of Science, University of AL- Mustansiriyah, Iraq <br> E-mail: lionwight_2009@yahoo.com 


#### Abstract

This paper concerned with study of vartional iteration method(VIM) and fractional variational iteration method (FVIM) to solve time fractional differential equations (FDE's). So FVIM is more effective than VIM in solving the FDE's. Keywords: PDE's; FDE's; fractional calculus; Riemann-Liouville; Lagrange multiplier.


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## 1. Introduction

The class of fraction calculus is one of the most convenient classes of (FDEs) which viewed as generalized differential equation $[1,2]$. In the sense that, much of the theory and, hence, applications of differential equation can be extended smoothly to (FDEs) with the same flavor and spirit of the realm of differential equation. (FDEs) have been proved to be a valuable tool in modeling many phenomena in the fields of physics, chemistry, engineering, aerodynamics, electrodynamics of complex medium, polymer rheology, and so forth [17-20,7,8,22,26,28] . The (VIM), which was first proposed by He et al. [9-14] and has been shown to be very efficient for handling a wide class of physical problems. As early as 1998, the variational iteration method was shown to be an effective tool for factional calculus [15]; hereafter, the method has been routinely used to solve various (FDEs). [23,16,3-5,21,29] for many years ,the (VIM and FVIM) is relatively new and effective approaches to find the approximate solution of PDEs, because they provide immediate and visible symbolic terms of analytic solutions, as well as numerical approximate solutions to both linear and nonlinear (PDEs) without linearization or discretization.

## 2. Preliminaries and Notations

In this section, we describe some necessary definitions and mathematical preliminaries of the fractional calculus theory.
Definition 1, [25]: A real function $h(t), t>0$, is said to be in the space $C_{\mu}, \mu \in R$, if there exists a real number $p>\mu$, such that $h(t)=t p h_{1}(t)$, where $h_{1}(t) \in C(0, \infty)$, and it is said to be in the space $C^{n}{ }_{\mu}$ if and only if $h^{(n)} \in C_{\mu}, n \in N$.
Definition 2 ,[24]. Riemann-Liouville fractional integral operator $\left(J^{\alpha}\right)$ of order $\alpha \geq 0$, of a function $f \in C_{\mu}, \mu \geq-1$ is defined as
$J^{\alpha} f(t)=\frac{1}{\Gamma(\alpha)} \int_{0}^{t}(t-\tau)^{\alpha-1} f(\tau) d \tau, \mathrm{t}>0$
$J^{0} f(t)=f(t)$
$\Gamma(\alpha)$ is the well-known gamma function. Some properties of the operator $J^{\alpha}$ can be found in [6]. We give in the following for $f \in C_{\mu}, \mu \geq-1, \alpha, \beta \geq 0$ and $\gamma>-1$ :

1. $J^{\alpha} J^{\beta} f(t)=J^{\alpha^{+} \beta} f(t)$,
2. $J^{\alpha} J^{\beta} f(t)=J^{\beta} J^{\alpha} f(t)$,
3. $J^{\alpha} t^{\gamma}=\frac{\Gamma(\gamma+1)}{\Gamma(\alpha+\gamma+1)} x^{\alpha+\gamma}$.

The Riemann-Liouville derivative has certain disadvantages when trying to model real-world phenomena with FDEs . Therefore, we will introduce a modified fractional differential operator $D_{x}^{\alpha}$ proposed by Caputo [27].
Definition 3, [25]. The fractional derivative of $f(x)$ in the Caputosense is defined as

$$
\left(D_{x}^{\alpha} f\right)(\mathrm{x})=\left\{\begin{array}{l}
\frac{1}{\Gamma(m-\alpha)} \int_{0}^{x} \frac{f^{m}(\xi)}{(x-\xi)^{\alpha-m+1}} d \xi,(\alpha>0, m-1<\alpha<m)  \tag{3}\\
\frac{\partial^{m} f(x)}{\partial x^{m}}, \alpha=m
\end{array}\right.
$$

where $f: R \rightarrow R, x \rightarrow f(x)$ denotes a continuous (but not necessarily differentiable) function. Some useful formulas and results of modified Riemann- Liouville derivative, which we need here, are listed as follows:

$$
\begin{align*}
& D_{x}^{\alpha} c=0, \alpha>0, c=\text { constant } \\
& D_{x}^{\alpha}[c f(x)]=c D_{x}^{\alpha} f(x), \alpha>0, c=\text { constant } \tag{4}
\end{align*}
$$

$$
\begin{aligned}
D_{x}^{\alpha} x \beta & =\frac{\Gamma(1+\beta)}{\Gamma(1+\beta-\alpha)} x^{\beta-\alpha}, \beta>\alpha>0 \\
D_{x}^{\alpha}[f(x) g(x)] & =\left[D_{x}^{\alpha} f(x)\right] g(x)+f(x)\left[D_{x}^{\alpha} g(x)\right], \\
D_{x}^{\alpha}[f(x(t))] & =f_{x}^{\prime}(x) x^{(\alpha)}(t) .
\end{aligned}
$$

## 3. Variational Iteration Method, [25].

In this section, the (VIM) is introduced. Here a description of the method of fractional differential initial value problem:

$$
\begin{equation*}
D_{t}^{\alpha} u(x, t)+N[u(x, t)]+L[u(x, t)]=g(x, t), t>0 \tag{5}
\end{equation*}
$$

Where $L$ is the linear operater, $N$ is nonlinear operator in $x, t$, and $D^{\alpha}$ is modified Riemann-Liouville derivative of order $\alpha$, subject to the initial conditions.
$u^{(k)}(x, 0)=c_{k}(x), k=0,1,2, \ldots, m-1, m-1<\alpha<m$.
According to He's variational iteration method [9-14] from(2), we can construct a correction functional as follows:

$$
\begin{equation*}
u_{n+1}(t)=u_{n+1}(t)+\int_{0}^{t} \lambda(\tau)\left\{L\left(u_{n}(\tau)\right)+N\left(\tilde{u}_{n}(\tau)\right)-g(\tau)\right\}, n \geq 0 \tag{7}
\end{equation*}
$$

where $\lambda$ is a general Lagrange multiplier which can be optimally identified via variational theory and $\tilde{u}_{n}$ is a restricted variation which means $\delta \tilde{u}_{n}=0$, then several approximations $u_{\mathrm{n}}(t), n \geq 0$ follow immediately. Consequently, the exact solution may be obtained as:

$$
\begin{equation*}
u(t)=\lim _{n \rightarrow \infty} u_{n}(t) \tag{8}
\end{equation*}
$$

## 4. Fractional Variational Iteration Method,[25]:

We can construct a correction functional for (5) as follows:

$$
\begin{equation*}
u_{k+1}(x, t)=u_{k+1}(x, t)+\int_{0}^{t} \lambda(t, \tau)\left(D_{t}^{\alpha} u_{k}(x, \tau)+N\left[\tilde{u}_{k}(x, \tau)\right]+L\left[\tilde{u}_{k}(x, \tau)\right]-g(x, \tau)\right) d \tau \tag{9}
\end{equation*}
$$

Where $\tilde{u}(x, t)$ is a restricted variation. Taking Laplace transform to both sides of (9) as

$$
\begin{equation*}
\bar{u}_{k+1}(x, t)=\bar{u}_{k}(x, t)+\bar{L}\left[\int_{0}^{t} \lambda(t, \tau)\left(D_{t}^{\alpha} u_{k}(x, \tau)+N\left[\tilde{u}_{k}(x, \tau)\right]+L\left[\tilde{u}_{k}(x, \tau)\right]-g(x, \tau)\right) d \tau\right] \tag{10}
\end{equation*}
$$

Where $\bar{u}_{k}(x, t)$ is Laplace transform of $u_{k}(x, t)$ with respect to $t$ and $L$ is operator of Laplace transform. By assuming that the Lagrange multiplier has the form as
$\lambda(t, \tau)=\lambda(t-\tau)$, so that $\bar{L}\left[J_{\tau}^{\alpha} \lambda\left(D_{\tau}^{\alpha} u_{k}(x, \tau)+N\left[\tilde{u}_{k}(x, \tau)\right]+L\left[\tilde{u}_{k}(x, \tau)\right]-g(x, \tau)\right)\right]$.
Is the convolution of the function $\lambda(t)$ and $D_{t}^{\alpha} u_{k}(x, t)+N\left[\tilde{u}_{k}(x, t)\right]+L\left[\tilde{u}_{k}(x, t)\right]-g(x, t)$.
Because $\tilde{u}(x, t)$ is a restricted variation, we have
$\delta \bar{L}\left[J_{t}^{\alpha} \lambda\left(N\left[\tilde{u}_{k}(x, t)\right]+L\left[\tilde{u}_{k}(x, t)\right]-g(x, t)\right)\right]=0$
Taking the variation derivative $\delta$ on the both sides of (10) ,we can derive

$$
\begin{align*}
\delta \bar{u}_{k+1}(x, t)=\delta \bar{u}_{k}(x, t)+\delta \bar{L}\left[J_{t}^{\alpha} \lambda\right. & \left.\left(N\left[\tilde{u}_{k}(x, t)\right]+L\left[\tilde{u}_{k}(x, t)\right]-g(x, t)\right)\right]=0 \\
& =\left(1+\bar{\lambda}(s) s^{\alpha}\right) \delta \bar{u}_{k}(x, s) . \tag{12}
\end{align*}
$$

If setting the coefficient of $\delta \bar{u}_{k}(x, s)$ to zero, we can get

$$
\begin{equation*}
\bar{\lambda}(s)=-\frac{1}{s^{\alpha}} \tag{13}
\end{equation*}
$$

And the Lagrange multiplier can be identified by using the invers Laplace transform

$$
\begin{equation*}
\lambda(t, \tau)=-\frac{(t-\tau)^{\alpha-1}}{\Gamma(\alpha)}=\frac{(-1)(\tau-t)^{\alpha-1}}{\Gamma(\alpha)} \tag{14}
\end{equation*}
$$

Substituting (14) into (10) and using the definition of Riemann-Liouville fractional integral operator, we get the iteration formula as follows:

$$
\begin{equation*}
u_{k+1}(x, t)=u_{k}(x, t)-J_{t}^{\alpha}\left(D_{t}^{\alpha} u_{k}(x, t)+N\left[u_{k}(x, t)\right]+L\left[u_{k}(x, t)\right]-g(x, t)\right) \tag{15}
\end{equation*}
$$

## 4. Applications and Results

In this section, we will solve one example for performing comparative studies. The exact solution of the example is known for special case $\alpha=1$ and has been solved by using VIM, FVIM.

## Example (1) :

consider the initial value problem of the fractional differential equation :

$$
\begin{equation*}
D_{t}^{\alpha} u(x, t)+\frac{1}{2}\left(u^{2}\right)_{x}(x, t)-u_{x x}(x, t)=x+x t^{2}-2 t^{2}, t>0, x \in R, 0<\alpha \leq 1 \tag{16}
\end{equation*}
$$

With initial condition

$$
\begin{equation*}
u(x, 0)=0 \tag{17}
\end{equation*}
$$

By using the (VIM) is given by:

$$
\begin{align*}
& u_{k+1}(x, t)=u_{k}(x, t)-\int_{0}^{t}\left(\frac{\partial^{\alpha}}{\partial \xi^{\alpha}} u_{k}(x, \xi)+\frac{1}{2} \frac{\partial}{\partial x}\left(u_{k}(x, \xi)\right)^{2}-\frac{\partial^{2}}{\partial x^{2}}\left(u_{k}(x, \xi)-\left(x+x \xi^{2}-2 \xi^{2}\right) d \xi\right.\right. \\
& u_{0}(x, t)=0  \tag{18}\\
& u_{1}(x, t)=x\left(t+\frac{t^{3}}{3}\right)-2 \frac{t^{3}}{3} \\
& u_{2}(x, t)=x\left(2 t+\frac{t^{3}}{3}-2 \frac{t^{5}}{15}-\frac{t^{7}}{63}-\frac{t^{2-\alpha}}{\Gamma(3-\alpha)}+2 \frac{t^{4-\alpha}}{\Gamma(5-\alpha)}-4 \frac{t^{3}}{3}+2 \frac{t^{5}}{15}+2 \frac{t^{7}}{63}\right)
\end{align*}
$$

$$
u_{3}(x, t)=x\left(3 t-2 \frac{t^{3}}{3}-2 \frac{t^{5}}{15}-18 \frac{t^{7}}{63}-8 \frac{t^{11}}{2475}-19 \frac{t^{15}}{3969}-3 \frac{t^{2-\alpha}}{\Gamma(3-\alpha)}-\frac{t^{(3-\alpha)^{2}}}{(\Gamma(4-\alpha))^{2}}\right)
$$

$$
-20 \frac{t^{(5-\alpha)^{2}}}{(\Gamma(6-\alpha))^{2}}+16 \frac{t^{4-\alpha}}{3 \Gamma(5-\alpha)}+238 \frac{t^{6-\alpha}}{15 \Gamma(7-\alpha)}+5038 \frac{t^{8-\alpha}}{63 \Gamma(9-\alpha)}+\frac{t^{3-2 \alpha}}{\Gamma(4-2 \alpha)}
$$

$$
-2 \frac{t^{5-2 \alpha}}{\Gamma(6-2 \alpha)}-6 \frac{t^{3}}{3}+2 \frac{t^{5}}{15}+2 \frac{t^{7}}{63}
$$

$$
\begin{aligned}
u_{4}(x, t) & =4 x t-10 x \frac{t^{3}}{3}-2 x \frac{t^{5}}{15}-22 x \frac{t^{7}}{63}-12 x \frac{t^{11}}{2475}-9739 x \frac{t^{15}}{3969}-6 x \frac{\text { SSAS }}{\Gamma(3-\alpha)}+34 \frac{7-t^{4} @ 21}{\Gamma(5-\alpha)} \\
& +16 x \frac{t^{6-\alpha}}{\Gamma(7-\alpha)}+1440 x \frac{t^{8-\alpha}}{\Gamma(9-\alpha)}+129024 x \frac{t^{12-\alpha}}{\Gamma(13-\alpha)}+2 x \frac{t^{16-\alpha}}{\Gamma(17-\alpha)}+3 x \frac{t^{3-2 \alpha}}{\Gamma(4-2 \alpha)} \\
& -x \frac{\left(t^{3-\alpha}\right)^{2}}{(\Gamma(4-\alpha))^{2}}-18 x \frac{\left(t^{3-\alpha}\right)^{2}}{(\Gamma(4-\alpha))^{2}}-2 x \frac{\left(t^{4-\alpha}\right)^{2}}{(\Gamma(5-\alpha))^{4}}+x \frac{\left(t^{4-2 \alpha}\right)^{2}}{(\Gamma(5-2 \alpha))^{2}}-128 x \frac{t^{23}}{140889375} \\
& -722 x \frac{t^{31}}{488341791}+64 \frac{t^{4-\alpha}}{3 \Gamma(5-\alpha)}-2 \frac{t^{6-\alpha}}{15 \Gamma(7-\alpha)}-5042 \frac{t^{8-\alpha}}{63 \Gamma(9-\alpha)}-20 \frac{\left(t^{5-\alpha}\right)^{2}}{(\Gamma(6-\alpha))^{2}} \\
& +\frac{t^{3-2 \alpha}}{\Gamma(4-2 \alpha)}-22 \frac{t^{5-2 \alpha}}{3 \Gamma(6-2 \alpha)}-238 \frac{t^{7-2 \alpha}}{15 \Gamma(8-2 \alpha)}-5038 \frac{t^{9-2 \alpha}}{63 \Gamma(10-2 \alpha)}+20 \frac{\left(t^{6-2 \alpha}\right)^{2}}{(\Gamma(7-2 \alpha))^{2}} \\
& -\frac{t^{4-3 \alpha}}{\Gamma(5-3 \alpha)}+2 \frac{t^{6-3}}{\Gamma(7-3 \alpha)}-8 \frac{t^{3}}{3}+2 \frac{t^{5}}{15}+2 \frac{t^{7}}{63}
\end{aligned}
$$

Table (1) . Numerical values by using VIM when $\alpha=0.9$

| T | X | $\mathbf{u i n}_{\mathrm{i}}$ | Exact solution | Absolute error |
| :---: | :---: | :---: | :---: | :---: |
| 0.01 | 0 | 0 | 0 | 0 |
|  | 0.1 | 0.0009993 | 0.001 | 13E-7 |
|  | 0.2 | 0.0019994 | 0.002 | 6E-7 |
|  | 0.3 | 0.0029996 | 0.003 | 5E-7 |
|  | 0.4 | 0.0039997 | 0.004 | 53E-7 |
|  | 0.5 | 0.0049995 | 0.005 | 5E-7 |
|  | 0.6 | 0.0059995 | 0.006 | 46E-7 |
|  | 0.7 | 0.0069996 | 0.007 | 43E-7 |
|  | 0.8 | 0.0079996 | 0.008 | 4E-7 |
|  | 0.9 | 0.0089996 | 0.009 | 33E-7 |
|  | 1 | 0.0099996 | 0.01 | 3E-7 |
| 0.02 | 0 | 0 | 0 | 0 |
|  | 0.1 | 0.0019949 | 0.002 | 50E-7 |
|  | 0.2 | 0.0039951 | 0.004 | 48E-7 |
|  | 0.3 | 0.0059954 | 0.006 | 45E-7 |
|  | 0.4 | 0.0079957 | 0.008 | 42E-7 |
|  | 0.5 | 0.009996 | 0.01 | 4E-6 |
|  | 0.6 | 0.0119962 | 0.012 | 37E-7 |
|  | 0.7 | 0.0139965 | 0.014 | 34E-7 |
|  | 0.8 | 0.0159968 | 0.016 | 32E-7 |
|  | 0.9 | 0.0179971 | 0.018 | 29E-7 |
|  | 1 | 0.0199973 | 0.02 | 26E-7 |

Table (1) shows the approximate solution for (16) obtained for value $\alpha$ using VIM . the values of $\alpha=1$ is the only case for which know the exact solution $u(x, t)=x t$.


Figure(1): Exact solution of example (1)


Figure(2): Approximate solution of example (1) by VIM
Now, to solve problem (16) by (FVIM) , we can obtain the following approximation:
$u_{k+1}(x, t)=u_{k}(x, 0)+J^{\alpha}\left(x+x t^{2}-2 t^{2}\right)-J^{\alpha}\left(\frac{1}{2}\left(u_{k}{ }^{2}\right)_{x}-\left(u_{k}\right)_{x x}\right)$
we get:
$u_{0}(x, t)=0$
$u_{1}(x, t)=x\left(\frac{t^{\alpha}}{\Gamma(1+\alpha)}+\frac{2 t^{2+\alpha}}{\Gamma(3+\alpha)}\right)-\frac{4 t^{2+\alpha}}{\Gamma(3+\alpha)}$

$$
\begin{aligned}
u_{3}(x, t)= & x\left(\frac{t^{\alpha}}{\Gamma(1+\alpha)}+\frac{2 t^{2+\alpha}}{\Gamma(3+\alpha)}-\frac{t^{3+\alpha}}{(\Gamma(1+\alpha))^{2}(\Gamma(3 \alpha+1))}-\frac{4 t^{4+3 \alpha}}{(\Gamma(3+\alpha))^{2}(\Gamma(3 \alpha+5))}\right. \\
& -\frac{4 t^{2+\alpha}}{\Gamma(3+\alpha)} \\
u_{3}(x, t)= & x\left(\frac{t^{\alpha}}{\Gamma(1+\alpha)}+\frac{2 t^{2+\alpha}}{\Gamma(3+\alpha)}-\frac{2 t^{3 \alpha}}{(\Gamma(1+\alpha))^{2}(\Gamma(3 \alpha+1))}-\frac{8 t^{4+3 \alpha}}{(\Gamma(3+\alpha))^{2}(\Gamma(3 \alpha+5)}\right. \\
& \left.+\frac{t^{7 \alpha}}{(\Gamma(1+\alpha))^{4}(\Gamma(3 \alpha+1))^{2}(\Gamma(7 \alpha+1)}-\frac{16 t^{8+7 \alpha}}{(\Gamma(3+\alpha))^{4}(\Gamma(3 \alpha+5))^{2}(\Gamma(9+7 \alpha)}\right) \\
& -\frac{4 t^{2+\alpha}}{\Gamma(3+\alpha)} \\
u_{4}(x, t)= & x\left(\frac{t^{\alpha}}{\Gamma(1+\alpha)}+\frac{2 t^{2+\alpha}}{\Gamma(3+\alpha)}-\frac{3 t^{3 \alpha}}{(\Gamma(1+\alpha))^{2}(\Gamma(3 \alpha+1))}\right. \\
& -\frac{12 t^{4+3 \alpha}}{(\Gamma(3+\alpha))^{2}(\Gamma(3 \alpha+5)}+\frac{3 t^{7 \alpha}}{(\Gamma(1+\alpha))^{4}(\Gamma(3 \alpha+1))^{2}(\Gamma(7 \alpha+1)} \\
& -\frac{32 t^{8+7 \alpha}}{(\Gamma(3+\alpha))^{4}(\Gamma(3 \alpha+5))^{2}(\Gamma(7 \alpha+9)}-\frac{256 t^{16+14 \alpha}}{(\Gamma(1+\alpha))^{8}(\Gamma(3 \alpha+1)(\Gamma(7 \alpha+1)(\Gamma(14 \alpha+1)} \\
& -\frac{t^{14 \alpha}}{(\Gamma(3+\alpha))^{8}\left(\Gamma ( 3 \alpha + 5 ) ^ { 4 } \left(\Gamma(7 \alpha+9)^{2}(\Gamma(14 \alpha+17)\right.\right.}-\frac{4 t^{2+\alpha}}{\Gamma(3+\alpha)}
\end{aligned}
$$

## Table -(2) Numerical Solution by using FVIM when $\alpha=0.9$



| 0.3 | 0.0088680 | 0.006 | 2868E-7 |
| :---: | :---: | :---: | :---: |
| 0.4 | 0.0118257 | 0.008 | 38257E-7 |
| 0.5 | 0.0147835 | 0.01 | 47835E-7 |
| 0.6 | 0.0177413 | 0.012 | 57413E-7 |
| 0.7 | 0.0206991 | 0.014 | $66991 \mathrm{E}-7$ |
| 0.8 | 0.0236569 | 0.016 | $76569 \mathrm{E}-7$ |
| 0.9 | 0.0266147 | 0.018 | $86147 \mathrm{E}-7$ |
| 1 | 0.0295724 | 0.02 | $95724 \mathrm{E}-7$ |

Table (2) shows the approximate solution for (16) obtained for value $\alpha$ using methods FVIM. The values of $\alpha=1$ is the only case for which know the exact solution $u(x, t)=x t$.


Figure (3): Exact solution of example (1)


Figure(4): Approximate solution of example (1) by FVIM

## 5. Conclusion:

The VIM has been successfully employed to obtain the approximate solution of FDE's, FVIM is powerful and effective tools for solution of FDE's and give approximation of higher accuracy, reduces the computational workload by avoiding the evaluation of VIM. So FVIM is more effective than VIM in solving the FDE's. The compared results of (VIM and FVIM) to solve FDE with the exact solutions also give a good comparative between two method, using Programming MATLAB to get the results of this method.

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